Quantum Phase Transitions in Magnetic Impurity Problems

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Project Goals

Physics Goals: Explore quantum phase transitions in strongly correlated electron systems in order to

- elucidate non-Fermi-liquid properties;
- understand "local criticality" in heavy fermions.

IT Goals: Advance the numerical RG method by providing

- new algorithms (e.g., for Bose-Fermi problems);
- efficient, adaptable codes for "complex" impurity problems (e.g., orbital degeneracy, coupled quantum dots, DMFT & beyond).

"Classic" Quantum Impurity Models

- Describe local, dynamical degree of freedom coupled to dispersive bath(s) of noninteracting (quasi)particles.
- ► E.g., Kondo model for a spin S coupled to a conduction band $H = J \mathbf{S} \cdot \mathbf{s} + H_{cond},$

where

$$\mathbf{s} = \frac{1}{2} \sum_{\sigma,\sigma'} c_{0\sigma}^{\dagger} \boldsymbol{\sigma}_{\sigma\sigma'} c_{0\sigma'}, \quad H_{\text{cond}} = \sum_{\mathbf{k},\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^{\dagger} c_{\mathbf{k}\sigma}, \quad |\epsilon_{\mathbf{k}}| \le D.$$

- Breakdown of perturbative expansion in J/D prompted development of many techniques.
 - Approximate: perturbative scaling/RG, large-degeneracy.
 - Exact (but limited): Bethe ansatz, bosonization.
 - Numerical (controlled): numerical RG, Quantum Monte Carlo.

"Modern" Quantum Impurity Models

- ► Feature some or all of the following:
 - complex impurities with many internal degrees of freedom;
 - multiple impurities and/or multiple conduction bands;
 - coupling to bosonic baths.
- ► Topical examples include . . .
 - ▷ Magnetic clusters on metallic surfaces [Crommie; Manoharan; . . .]
 - Coupled quantum dots
 - Cluster corrections to the dynamical mean-field theory of correlated lattice fermions [Jarrell et al. (1998); Kotliar et al. (2001)]
 - Bose-Fermi impurity models:
 - Enter the extended dynamical mean-field treatment of heavy fermions [Si et al. (2001)].
 - Describe quantum dots coupled to noisy leads [Le Hur (2004)].

Critical Local Moments in Lattice Problems

The Bose-Fermi Kondo model arises in the extended DMFT treatment of the Kondo lattice [Si et al. (2001, 2003)]:

- Extended DMFT includes some spatial fluctuations [Smith & Si (2000); Chitra & Kotliar (2000)].
- Effective impurity model couples a
 local spin to a
 fermionic bath
 and a vector bosonic bath.
- Bath densities of states must be determined self-consistently.



EDMFT Prediction: Two Types of Quantum Criticality

conventional criticality

"local criticality"



only long-wavelength fluctuations are important



long-wavelength and spatially local (dynamical) fluctuations play central roles

The locally critical QCP reproduces some anomalous properties of $CeCu_{6-x}Au_x$ and $YbRh_2Si_2$.

The Bose-Fermi Kondo Model

- Describes a local spin S coupled both to delocalized fermions (e.g., a conduction band) and bosons (e.g., phonons).
- Isotropic version has a Hamiltonian

$$H = \mathbf{J} \mathbf{S} \cdot \mathbf{s} + H_{\text{Fermi}} + \mathbf{g} \mathbf{S} \cdot \mathbf{u} + H_{\text{Bose}},$$

where (for
$$\alpha = x$$
, y , z)

$$\begin{split} s_{\alpha} &= \frac{1}{2} \sum_{\sigma,\sigma'} c_{0\sigma}^{\dagger} \,\sigma_{\sigma\sigma'}^{\alpha} \,c_{0\sigma'} & u_{\alpha} &= a_{0\alpha} + a_{0\alpha}^{\dagger}, \\ H_{\text{Fermi}} &= \sum_{\mathbf{k},\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^{\dagger} c_{\mathbf{k}\sigma} & H_{\text{Bose}} &= \sum_{\mathbf{q},\alpha} \omega_{\mathbf{q}} a_{\mathbf{q}\alpha}^{\dagger} a_{\mathbf{q}\alpha} \end{split}$$

Anisotropic versions distinguish between

$$egin{array}{rcl} J_z & ext{and} & J_x = J_y = J_ot, \ g_z & ext{and} & g_x = g_y = g_ot. \end{array}$$

The Bose-Fermi Kondo Model (continued)

 $H = \mathbf{J} \mathbf{S} \cdot \mathbf{s} + H_{\text{Fermi}} + \mathbf{g} \mathbf{S} \cdot \mathbf{u} + H_{\text{Bose}}.$

Take a flat fermionic density of states:



Perturbative Solutions of the Bose-Fermi Kondo Model

- ► The model has been solved by expansion in \(\epsilon = 1 s\) [Si & Smith (1999), Sengupta (2000), Zhu & Si (2002), Zaránd & Demler (2002)].
- A QCP separates Kondo and bosonic regimes.
- Critical point couplings $\rho_0 J^*$ and $K_0 g^*$ are of order ϵ (except for $g_{\perp} = 0$).
- At QCP, χ_{loc} shows power laws in ω and T with εdependent exponents.



• Locally critical EDMFT solution corresponds to $\epsilon = 1^-$. \Rightarrow Seek non-perturbative solutions.

Numerical Renormalization Group Method [Wilson (1974)]

- ► Replaces a continuum of fermionic states (|ε| ≤ D) by a discrete set having energies ε = ±D, ±DΛ⁻¹, ±DΛ⁻², ... (Λ > 1).
- ► Then the kinetic energy is converted to a tight-binding form:

$$H_{\text{Fermi}} = D \sum_{\sigma} \sum_{n=0}^{\infty} \Lambda^{-n/2} \left(c_{n,\sigma}^{\dagger} c_{n-1,\sigma}^{\dagger} + \text{h.c.} \right),$$

where only $c_{0,\sigma}$ couples to the impurity.



The exponential decay of the hopping permits iterative solution via diagonalization of progressively longer chains.

Discretizing a Bosonic Bath

- Can use the same energy discretization as for fermions.
- No negative- ω states \Rightarrow hopping coefficients decay faster:

$$H_{\text{Bose}} = D \sum_{\alpha} \sum_{n=0}^{\infty} \left[t \Lambda^{-n} \left(a_{n,\alpha}^{\dagger} a_{n-1,\alpha} + \text{h.c.} \right) + e \Lambda^{-n} a_{n,\alpha}^{\dagger} a_{n,\alpha} \right]$$



This discretization has been used to study the spin-boson model [Bulla et al. (2003)].

Combining Fermionic and Bosonic Baths

- Seek an iterative procedure that treats simultaneously fermionic and bosonic degrees of freedom having the same energy scale.
- One method: add a bosonic site at every second iteration:



Iterate until reach a scale-invariant fixed point describing the ground state.

Adding Bosons = More CPU Time!

- Add a fermionic site \Rightarrow basis increases by a factor $N_F = 4$.
- ▶ Number of bosons on each site is unlimited, but restrict it to no more than N_b per bath. $4 \le N_b \le 8$ seems to suffice.
- ► Add both fermionic and bosonic sites \Rightarrow basis increases by $N_F = 4(N_b+1)^n$ for n Bose baths.
- ► After a few iterations, basis is so large that can retain only the M states of lowest-energy. Typically, 500 ≤ M ≤ 2000.
- ► As in other "modern" impurity problems, large N_F and M lead to long CPU times ~ O [(N_FM)³].
- To tackle these problems, we are developing parallelized NRG codes (using MPI, ScaLAPACK).

Initial Results: Bose-Fermi Kondo and Anderson Models

• Have started with the Ising-symmetry ($g_{\perp} = 0$) model:

$$H_{\rm imp} = J\mathbf{S} \cdot \mathbf{s} + gS_z(a_0 + a_0^{\dagger}).$$

Has smallest N_F , and is possibly the most relevant for CeCu_{6-x}Au_x.

- Have calculated phase diagram, T = 0 impurity dynamics, and static magnetic response.
- Will show results along constant-*J* cuts through the parameter space.





Increasing $\Delta < \Delta_c$ for s = 0.7



Kondo regime: A(0) is pinned, but the Abrikosov-Suhl resonance narrows as $\Delta \rightarrow \Delta_c$, signaling suppression of the Kondo effect.

Impurity Spectral Function [$\Delta = (K_0 g)^2$]

Decreasing $\Delta > \Delta_{\rm c}$ for s = 0.7



Bosonic regime: A low-energy feature grows as $\Delta \rightarrow \Delta_c$.

Impurity Spectral Function [$\Delta = (K_0 g)^2$]

s = 0.7



Passing through the quantum critical point, A(0) undergoes a jump.



Can extract the same critical exponent from the width of the Abrikosov-Suhl resonance and from the many-body eigenspectrum.

Static Susceptibility



Summary

- Have implemented the first numerical renormalization group treatment of a quantum impurity coupled to both fermionic and bosonic baths.
- As an initial application, are studying the Ising-symmetry Bose-Fermi Kondo model.
- The method should permit study of critical properties beyond the range of perturbative methods.
- The method will extend to other models and can serve as an impurity solver in extended DMFT treatments of lattice fermion problems.
- Parallelized, readily-adaptable NRG codes will be available for a wide range of quantum impurity problems.